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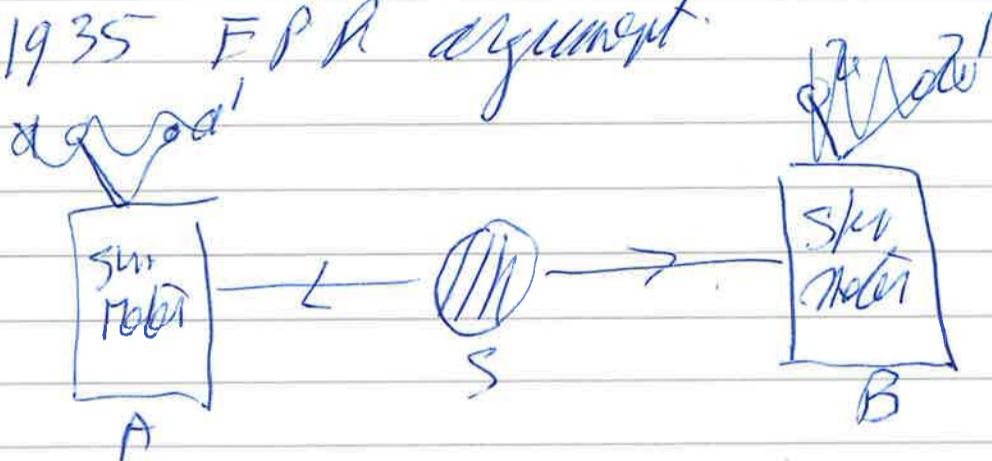
# Algebraic Proof of Nonlocality

Oxford February 1990

In what sense is QM a nonlocal theory?  
of 2 slit experiment, extended wavefunctions  
Efimov states in nuclear physics etc.

But R&FT is based on causality axioms,  
so how can the theory be nonlocal?

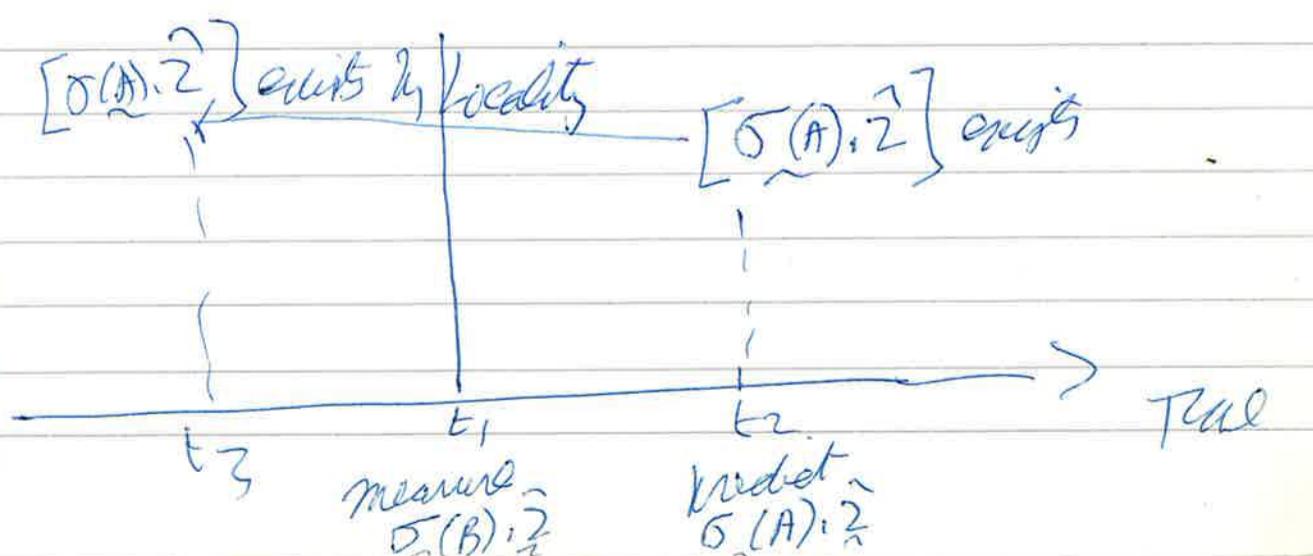
(1) 1935 EPR argument



set  $\hat{a} \parallel \hat{b} \parallel \hat{z}$  If  $S_{\text{spacelike}} = \frac{1}{2} \{ \langle \hat{S}(A) \cdot \hat{z} = +1 \rangle$

$|\hat{S}(B) \cdot \hat{z} = -1 \rangle - |\hat{S}(A) \cdot \hat{z} = -1 \rangle$

$|\hat{S}(B) \cdot \hat{z} = +1 \rangle \}$



(2)

for Everett's Diamond

$F \rightarrow \sim(L) \vee$  Incompleteness

C  $\mathcal{DM} \Rightarrow$  nonlocality or Incompleteness

↓  
Completed version  
of  $\mathcal{DM}$   
(hidden variables)

↓ Bell Inequalities  
violated by soft.  
nonlocality.

So  $\mathcal{DM}$  is nonlocal & incomplete.

But Bohr denied the nonlocality as a  
"physical" effect.  
So we reject Bohr's approach

then what about

? + Locality + hidden variables  $\Rightarrow$  Bell Th.

Early proof Bell (1964) against determinism  
+ probability  $\rightarrow$  statistics  
(S.D. for incompatible observables) + No necessary  
experiments  
cp. to -Shimony, 2011

Two Controversies (i) Does proof of Bell, under  
determinism, count as to T.D.?

(after Stapp 1971, 1977) Fano says yes (1982)  
Redhead says no (1983, 1988)  
with Svetlichny, Batterfield, Bentsen

(3)

Fins 1982 theorem extended to  
beautiful mathematical theory by Pitowsky  
and others - see Pitowsky (1981)  
Quantum Probability - Quantum Logic

Generalized Bell inequalities are fact  
inequalities defining sets of foams in  
a multidimensional polytope.

(2) Does Car the Staff - Standard proof  
to extended to indeterminism.

Staff says yes

Bellman Redhead (after, Butterfield) says no.

1982 Redhead (1983, 1987).

CBA paper 1990  
A Staff in the Way Pitowsky

Can one give proofs of nonlocality in  
h.v. mechanics that do not  
use probabilities? They?

Answer Assuming determinism yes (weakly)

Assuming indeterminism no

But they do not exist algebraic  
so-called stochastic algebraic proofs  
of nonlocality (EPR, EPRB)  
first placed by Hardy with general stochastic  
framework and places certain probabilities  
on the other way these assigned  
to get an algebraic contradiction.

## History of the algebraic field's

1.) Project: Define a Kochen-Specker (KS) configuration for two spin systems. (of Gleason (1957))

10. ~~say~~<sup>up</sup> local observables like  $S(A)$  ~~and~~<sup>must</sup> in a local context of properties of the whole system.

Inspired by Bell in 1976 in form of a question.

Caved Maczynski's Thesis (1971) he says that it is extendable from maximal to locally maximal observables?

But MS theorem was so extended by Hemmo, Humphreys and Bell in 1980.

So no algebraic proof of non-existence could be given.

2) In 1983 Redhead's Neuron defined 1 K-S paradox as one of a pair of spin-1/2 particles occurs separately and locally.  
Involved locally non-maximal observables.

3) Stairs (1983) followed by Braun Siegelberg (1990)

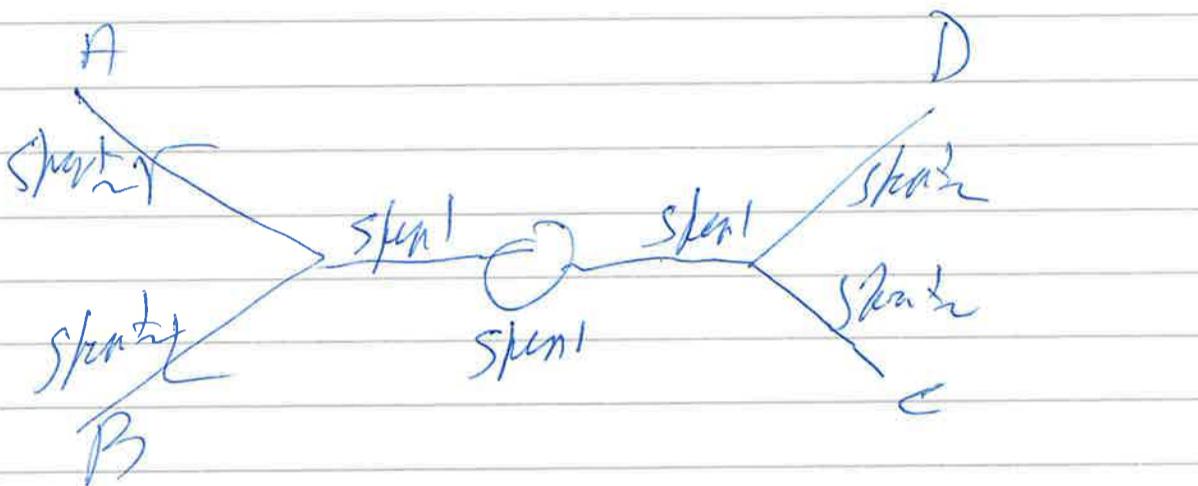
Forced law to get a proof of non-existence using similar structures to RST, but restricted to locally maximal observables but again other determinants.

(5)

4) Eby produced a moderate version of  
Span - Baus - Srofliat

5) Spangler, Horne, & Zelinger (1989)  
produced a broad new "algae"  
that, quite unrelated to K-S &  
leach.

6) 1990 Redhead's clinton showed the  
shift from green by Spangler was  
not due to shift to dark but  
clinton derived a moderate shift  
that does not.



Then it is shown that  $\text{out of } \theta_A + \theta_B - \theta_c - \theta_D = 11$   
 $\text{if } \theta_A + \theta_B - \theta_c - \theta_D = 0, A(\theta_A) B(\theta_B) C(\theta_c) D(\theta_D) = +1$   
 $A(\theta_A) B(\theta_B) C(\theta_c) D(\theta_D) = -1$

Consider the 5 possible settings for  $\{\theta_A, \theta_B, \theta_c, \theta_D\}$

	(1)	(2)	(3)	(4)	<u>(5)</u>	
$\theta_A \rightarrow$	$\downarrow$	$\rightarrow$	$\uparrow$	$\uparrow$	$\uparrow \downarrow$	$\left. \begin{array}{l} \text{From (4)(5)} \\ A(\uparrow) \\ = -A(\downarrow) \end{array} \right\}$
$\theta_B \rightarrow$	$\uparrow$	$\uparrow$	$\rightarrow$	$\uparrow$		
$\theta_c \rightarrow$	$\rightarrow$	$\uparrow$	$\uparrow$			
$\theta_D \rightarrow$	$\rightarrow$	$\rightarrow$	$\rightarrow$	$\rightarrow$		
$A B C D$	$= -1$	$= -1$	$= -1$	$= -1$	$= +1$	

6)

Now we obtain

$$\left. \begin{array}{l} (1) (2) \quad A(\rightarrow) B(\rightarrow) = A(\downarrow) B(\uparrow) \\ (3) \quad B(\rightarrow) C(\rightarrow) = B(\uparrow) C(\uparrow) \\ \text{similar} \quad A(\rightarrow) C(\rightarrow) = A(\uparrow) C(\uparrow) \end{array} \right\}$$

Similarly  $I = A(\downarrow) A(\uparrow)$

$\sim A(\downarrow) = A(\uparrow)$ , but  $A(\downarrow) = -A(\uparrow)$   
from (4) & (5)  
 $\therefore$  contradiction

Soergel considers this argument a counterargument to Birk's opposition to EPR. He fails to notice that discussion to the "superclassical case" where an element of reality exists by virtue of "perfect predictability".

But this is running into the question

$$A(\rightarrow), A(\uparrow) A(\downarrow) B(\rightarrow), B(\uparrow), C(\rightarrow), C(\uparrow), D(\rightarrow)$$

involved in the argument do not all exist in the same experimental context, or are they the products of different agents, and this struggle is decisive.

7) Soergel has generalized the EPR argument to the "classical case", in a manner parallel to that EPR did for Sklar

(7)

so where are we left?

conditions that cannot be explained

not causal

ep. polarities mostly Redlef (1986)  
and non-symmetry theories of Shiroki  
Kusano, Weber (1980) and others  
(including Gerhard, Page, Sherry, Redlef  
etc.)

Recent work on violations of Bell Suggestion  
in R&FT derived by T. Sumino, Werner,  
Leordan, Farber, etc. - also in press  
See also p. 115 of my book